

Remarks on $A_n^{(1)}$ face weights

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Abstract Elementary proofs are presented for the factorization of the elliptic Boltzmann weights of the $A_n^{(1)}$ face model, and for the sum-to-1 property in the trigonometric limit, at a special point of the spectral parameter. They generalize recent results obtained in the context of the corresponding trigonometric vertex model.

1 Introduction

In the recent work [8], the quantum R matrix for the symmetric tensor representation of the Drinfeld-Jimbo quantum affine algebra $U_q(A_n^{(1)})$ was revisited. A new factorized formula at a special value of the spectral parameter and a certain sum rule called sum-to-1 were established. These properties have led to vertex models that can be interpreted as integrable Markov processes on one-dimensional lattice including several examples studied earlier [7, Fig.1,2]. In this note we report analogous properties of the Boltzmann weights for yet another class of solvable lattice models known as IRF (interaction round face) models [2] or face models for short. More specifically, we consider the elliptic fusion $A_n^{(1)}$ face model corresponding to the symmetric tensor representation [6, 5]. For $n = 1$, it reduces to [1] and [4] when the fusion degree is 1 and general, respectively. There are restricted and unrestricted versions of the model. The trigonometric case of the latter reduces to the $U_q(A_n^{(1)})$ vertex model when the site variables tend to infinity. See Proposition 1. In this sense Theorem 1 and Theorem 2 given below, which are concerned with the unrestricted version, provide generalizations of [8, Th.2] and [8, eq.(30)] so as to include finite site variables (and also to the elliptic case in the former). In Section 3 we will also

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comment on the restricted version and difficulties to associate integrable stochastic models.

2 Results

Let $\theta_1(u) = \theta_1(u, p) = 2p^{\frac{1}{4}} \sin \pi u \prod_{k=1}^{\infty} (1 - 2p^{2k} \cos 2\pi u + p^{4k})(1 - p^{2k})$ be one of the Jacobi theta function ($|p| < 1$) enjoying the quasi-periodicity

$$\theta_1(u+1; e^{\pi i \tau}) = -\theta_1(u; e^{\pi i \tau}), \quad \theta_1(u+\tau; e^{\pi i \tau}) = -e^{-\pi i \tau - 2\pi i u} \theta_1(u; e^{\pi i \tau}), \quad (1)$$

where $\text{Im} \tau > 0$. We set

$$[u] = \theta_1\left(\frac{u}{L}, p\right), \quad [u]_k = [u][u-1] \cdots [u-k+1], \quad \begin{bmatrix} u \\ k \end{bmatrix} = \frac{[u]_k}{[k]_k} \quad (k \in \mathbb{Z}_{\geq 0}), \quad (2)$$

with a nonzero parameter L . These are elliptic analogue of the q -factorial and the q -binomial:

$$(z)_m = (z; q)_m = \prod_{i=0}^{m-1} (1 - zq^i), \quad \binom{m}{l}_q = \frac{(q)_m}{(q)_l (q)_{m-l}}.$$

For $\alpha = (\alpha_1, \dots, \alpha_k)$ with any k we write $|\alpha| = \alpha_1 + \dots + \alpha_k$. The relation $\beta \geq \gamma$ or equivalently $\gamma \leq \beta$ means $\beta_i \geq \gamma_i$ for all i .

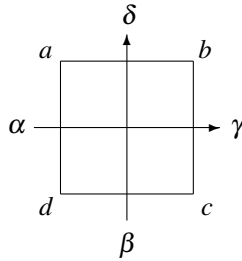
We take the set of local states as $\mathcal{P} = \eta + \mathbb{Z}^{n+1}$ with a generic $\eta \in \mathbb{C}^{n+1}$. Given positive integers l and m , let $a, b, c, d \in \mathcal{P}$ be the elements such that

$$\alpha = d - a \in B_l, \quad \beta = c - d \in B_m, \quad \gamma = c - b \in B_l, \quad \delta = b - a \in B_m, \quad (3)$$

where B_m is defined by

$$B_m = \{\alpha = (\alpha_1, \dots, \alpha_{n+1}) \in \mathbb{Z}_{\geq 0}^{n+1} \mid |\alpha| = m\}. \quad (4)$$

The relations (3) imply $\alpha + \beta = \gamma + \delta$. The situation is summarized as



To the above configuration round a face we assign a function of the spectral parameter u called Boltzmann weight. Its unnormalized version, denoted by $\bar{W}_{l,m} \left(\begin{smallmatrix} a & b \\ d & c \end{smallmatrix} \middle| u \right)$, is constructed from the $l = 1$ case as follows:

$$\begin{aligned} \bar{W}_{l,m} \left(\begin{smallmatrix} a & b \\ d & c \end{smallmatrix} \middle| u \right) &= \sum \prod_{i=0}^{l-1} \bar{W}_{1,m} \left(\begin{smallmatrix} a^{(i)} & b^{(i)} \\ a^{(i+1)} & b^{(i+1)} \end{smallmatrix} \middle| u - i \right), \\ \bar{W}_{1,m} \left(\begin{smallmatrix} a & b \\ d & c \end{smallmatrix} \middle| u \right) &= \frac{[u + b_\nu - a_\mu] \prod_{j=1}^{n+1} [b_\nu - a_j + 1]}{\prod_{j=1}^{n+1} [c_\nu - b_j]} \quad (d = a + \mathbf{e}_\mu, \quad c = b + \mathbf{e}_\nu), \end{aligned} \quad (5)$$

where $\mathbf{e}_i = (0, \dots, 0, \overset{i\text{th}}{1}, 0, \dots, 0)$. In (5), $a^{(0)}, \dots, a^{(l)} \in \tilde{\mathcal{P}}$ is a path from $a^{(0)} = a$ to $a^{(l)} = d$ such that $a^{(i+1)} - a^{(i)} \in B_1$ ($0 \leq i < l$). The sum is taken over $b^{(1)}, \dots, b^{(l-1)} \in \tilde{\mathcal{P}}$ satisfying the conditions $b^{(i+1)} - b^{(i)} \in B_1$ ($0 \leq i < l$) with $b^{(0)} = b$ and $b^{(l)} = c$. It is independent of the choice of $a^{(1)}, \dots, a^{(l-1)}$ (cf. [4, Fig.2.4]). We understand that $\bar{W}_{l,m} \left(\begin{smallmatrix} a & b \\ d & c \end{smallmatrix} \middle| u \right) = 0$ unless (3) is satisfied for some $\alpha, \beta, \gamma, \delta$.

The normalized weight is defined by

$$W_{l,m} \left(\begin{smallmatrix} a & b \\ d & c \end{smallmatrix} \middle| u \right) = \bar{W}_{l,m} \left(\begin{smallmatrix} a & b \\ d & c \end{smallmatrix} \middle| u \right) \frac{[1]_l^m}{[l]_l^m}. \quad (6)$$

It satisfies [5] the (unrestricted) star-triangle relation (or dynamical Yang-Baxter equation) [2]:

$$\begin{aligned} \sum_g W_{k,m} \left(\begin{smallmatrix} a & b \\ f & g \end{smallmatrix} \middle| u \right) W_{l,m} \left(\begin{smallmatrix} f & g \\ e & d \end{smallmatrix} \middle| v \right) W_{k,l} \left(\begin{smallmatrix} b & c \\ g & d \end{smallmatrix} \middle| u - v \right) \\ = \sum_g W_{k,l} \left(\begin{smallmatrix} a & g \\ f & e \end{smallmatrix} \middle| u - v \right) W_{l,m} \left(\begin{smallmatrix} a & b \\ g & c \end{smallmatrix} \middle| v \right) W_{k,m} \left(\begin{smallmatrix} g & c \\ e & d \end{smallmatrix} \middle| u \right), \end{aligned} \quad (7)$$

where the sum extends over $g \in \tilde{\mathcal{P}}$ giving nonzero weights. Under the same setting (3) as in (6), we introduce the product

$$S_{l,m} \left(\begin{smallmatrix} a & b \\ d & c \end{smallmatrix} \right) = [l]^{-1} \prod_{1 \leq i, j \leq n+1} \frac{[c_i - d_j]_{c_i - b_i}}{[c_i - b_j]_{c_i - b_i}}. \quad (8)$$

Note that $S_{l,m} \left(\begin{smallmatrix} a & b \\ d & c \end{smallmatrix} \right) = 0$ unless $d \leq b$ because of the factor $\prod_{i=1}^{n+1} [c_i - d_i]_{c_i - b_i}$. The following result giving an explicit factorized formula of the weight $W_{l,m}$ at special value of the spectral parameter is the elliptic face model analogue of [8, Th.2].

Theorem 1. *If $l \leq m$, the following equality is valid:*

$$W_{l,m} \left(\begin{smallmatrix} a & b \\ d & c \end{smallmatrix} \middle| u = 0 \right) = S_{l,m} \left(\begin{smallmatrix} a & b \\ d & c \end{smallmatrix} \right). \quad (9)$$

Proof. We are to show

$$\overline{W}_{l,m} \left(\begin{array}{c} a \ b \\ d \ c \end{array} \middle| 0 \right) = \frac{[l]_l}{[1]^l} \prod_{i,j} \frac{[c_i - d_j]_{c_i - b_i}}{[c_i - b_j]_{c_i - b_i}}. \quad (10)$$

Here and in what follows unless otherwise stated, the sums and products are taken always over $1, \dots, n+1$ under the condition (if any) written explicitly. We invoke the induction on l . It is straightforward to check (10) for $l = 1$. By the definition (5) the $l+1$ case is expressed as

$$\overline{W}_{l+1,m} \left(\begin{array}{c} a \ b \\ d \ c \end{array} \middle| 0 \right) = \sum_{\nu} \overline{W}_{l,m} \left(\begin{array}{c} a \ b \\ d' \ c' \end{array} \middle| 0 \right) \overline{W}_{1,m} \left(\begin{array}{c} d' \ c' \\ d \ c \end{array} \middle| -l \right) \quad (d' = d - \mathbf{e}_{\mu}, c' = c - \mathbf{e}_{\nu})$$

for some fixed $\mu \in [1, n+1]$. Due to the induction hypothesis on $\overline{W}_{l,m}$, the equality to be shown becomes

$$\begin{aligned} & \sum_{\nu} \frac{[l]_l}{[1]^l} \left(\prod_{i,j} \frac{[c'_i - d'_j]_{c'_i - b_i}}{[c'_i - b_j]_{c'_i - b_i}} \right) \frac{[-l + c'_{\nu} - d'_{\mu}] \prod_{k \neq \mu} [c'_{\nu} - d'_k + 1]}{\prod_k [c_{\nu} - c'_k]} \\ &= \frac{[l+1]_{l+1}}{[1]^{l+1}} \prod_{i,j} \frac{[c_i - d_j]_{c_i - b_i}}{[c_i - b_j]_{c_i - b_i}}. \end{aligned} \quad (11)$$

After removing common factors using $c'_i = c_i - \delta_{i\nu}$, $d'_i = d_i - \delta_{i\mu}$, one finds that (11) is equivalent to

$$\sum_{\nu} [c_{\nu} - d_{\mu} - l] \prod_{i \neq \nu} \frac{[c_i - d_{\mu} + 1]}{[c_{\nu} - c_i]} \prod_j [c_{\nu} - b_j] = [l+1] \prod_i [b_i - d_{\mu} + 1]$$

with l determined by $l+1 = \sum_j (c_j - b_j)$. One can eliminate d_{μ} and rescale the variables by $(b_j, c_j) \rightarrow (Lb_j + d_{\mu}, Lc_j + d_{\mu})$ for all j . The resulting equality follows from Lemma 1.

Lemma 1. *Let $b_1, \dots, b_n, c_1, \dots, c_n \in \mathbb{C}$ be generic and set $s = \sum_{i=1}^n (c_i - b_i)$. Then for any $n \in \mathbb{Z}_{\geq 1}$ the following identity holds:*

$$\sum_{i=1}^n \theta_1(z + c_i - s) \prod_{j=1, (j \neq i)}^n \frac{\theta_1(z + c_j)}{\theta_1(c_i - c_j)} \prod_{j=1}^n \theta_1(c_i - b_j) = \theta_1(s) \prod_{i=1}^n \theta_1(z + b_i).$$

Proof. Denote the LHS – RHS by $f(z)$. From (1) we see that $f(z)$ satisfies (12) with $B = \frac{n}{2}$, $A_1 = \frac{n(1+\tau)}{2} + \sum_{j=1}^n b_j$ and $A_2 = n$. Moreover it is easily checked that $f(z)$ possesses zeros at $z = -c_1, \dots, -c_n$. Therefore Lemma 2 claims $-(c_1 + \dots + c_n) - (B\tau + \frac{1}{2}A_2 - A_1) \equiv 0 \pmod{\mathbb{Z} + \mathbb{Z}\tau}$. But this gives $s \equiv 0$ which is a contradiction since b_j, c_j can be arbitrary. Therefore $f(z)$ must vanish identically.

Lemma 2. *Let $\text{Im}\tau > 0$. Suppose an entire function $f(z) \not\equiv 0$ satisfies the quasi-periodicity*

$$f(z+1) = e^{-2\pi i B} f(z), \quad f(z+\tau) = e^{-2\pi i(A_1+A_2z)} f(z). \quad (12)$$

Then $A_2 \in \mathbb{Z}_{\geq 0}$ holds and $f(z)$ has exactly A_2 zeros $z_1, \dots, z_{A_2} \pmod{\mathbb{Z} + \mathbb{Z}\tau}$. Moreover $z_1 + \dots + z_{A_2} \equiv B\tau + \frac{1}{2}A_2 - A_1 \pmod{\mathbb{Z} + \mathbb{Z}\tau}$ holds.

Proof. Let C be a period rectangle $(\xi, \xi+1, \xi+1+\tau, \xi+\tau)$ on which there is no zero of $f(z)$. From the Cauchy theorem the number of zeros of $f(z)$ in C is equal to $\int_C \frac{f'(z)}{f(z)} \frac{dz}{2\pi i}$. Calculating the integral by using (12) one gets A_2 . The latter assertion can be shown similarly by considering the integral $\int_C \frac{zf'(z)}{f(z)} \frac{dz}{2\pi i}$.

From Theorem 1 and (7) it follows that $S_{l,m} \begin{pmatrix} a & b \\ d & c \end{pmatrix}$ also satisfies the (unrestricted) star-triangle relation (7) without spectral parameter. The discrepancy of the factorizing points $u=0$ in (9) and “ $u=l-m$ ” in [8, Th.2] is merely due to a conventional difference in defining the face and the vertex weights.

Since (6) and (8) are homogeneous of degree 0 in the symbol $[\dots]$, the trigonometric limit $p \rightarrow 0$ may be understood as replacing (2) by $[u] = q^{u/2} - q^{-u/2}$ with generic $q = \exp \frac{2\pi i}{L}$. Under this prescription the elliptic binomial $\begin{bmatrix} m \\ l \end{bmatrix}$ from (2) is replaced by $q^{l(l-m)/2} \begin{pmatrix} m \\ l \end{pmatrix}_q$, therefore the trigonometric limit of (8) becomes

$$S_{l,m} \begin{pmatrix} a & b \\ d & c \end{pmatrix}_{\text{trig}} = \begin{pmatrix} m \\ l \end{pmatrix}_q^{-1} \prod_{1 \leq i, j \leq n+1} \frac{(q^{b_i-d_j+1})_{c_i-b_i}}{(q^{b_i-b_j+1})_{c_i-b_i}}. \quad (13)$$

The following result is a trigonometric face model analogue of [8, Th.6].

Theorem 2. *Suppose $l \leq m$. Then the sum-to-1 holds in the trigonometric case:*

$$\sum_b S_{l,m} \begin{pmatrix} a & b \\ d & c \end{pmatrix}_{\text{trig}} = 1, \quad (14)$$

where the sum runs over those b satisfying $c-d \in B_m$ and $d-a \in B_l$.

Proof. The relation (14) is equivalent to

$$\begin{pmatrix} m \\ l \end{pmatrix}_q = \sum_{\gamma \in B_l, \gamma \leq \beta} \prod_{1 \leq i, j \leq n+1} \frac{(q^{c_{ij}-\gamma_i+\beta_j+1})_{\gamma_i}}{(q^{c_{ij}-\gamma_i+\gamma_j+1})_{\gamma_i}} \quad (c_{ij} = c_i - c_j) \quad (15)$$

for any fixed $\beta = (\beta_1, \dots, \beta_{n+1}) \in B_m, l \leq m$ and the parameters c_1, \dots, c_{n+1} , where the sum is taken over $\gamma \in B_l$ (4) under the constraint $\gamma \leq \beta$. In fact we are going to show

$$\frac{(w_1^{-1} \dots w_n^{-1} q^{-l+1})_l}{(q)_l} = \sum_{|\gamma|=l} \prod_{1 \leq i, j \leq n} \frac{(q^{-\gamma_i+1} z_i / (z_j w_j))_{\gamma_i}}{(q^{\gamma_i-\gamma_i+1} z_i / z_j)_{\gamma_i}} \quad (l \in \mathbb{Z}_{\geq 0}), \quad (16)$$

where the sum is over $\gamma \in \mathbb{Z}_{\geq 0}^n$ such that $|\gamma| = l$, and $w_1, \dots, w_n, z_1, \dots, z_n$ are arbitrary parameters. The relation (15) is deduced from (16) $_{n \rightarrow n+1}$ by setting

$z_i = q^{c_i}, w_i = q^{-\beta_i}$ and specializing β_i 's to nonnegative integers. In particular, the constraint $\gamma \leq \beta$ automatically arises from the $i = j$ factor $\prod_{i=1}^n (q^{-\gamma_i + 1 + \beta_i})_{\gamma_i}$ in the numerator. To show (16) we rewrite it slightly as

$$q^{\frac{l^2}{2}} \frac{(w_1 \dots w_n)_l}{(q)_l} = \sum_{|\gamma|=l} \prod_{i=1}^n q^{\frac{\gamma_i^2}{2}} \frac{(w_i)_{\gamma_i}}{(q)_{\gamma_i}} \prod_{1 \leq i \neq j \leq n} \frac{(z_j w_j / z_i)_{\gamma_i}}{(q^{-\gamma_j} z_j / z_i)_{\gamma_i}}. \quad (17)$$

Denote the RHS by $F_n(w_1, \dots, w_n | z_1, \dots, z_n)$. We will suppress a part of the arguments when they are kept unchanged in the formulas. It is easy to see

$$F_n(w_1, w_2 | z_1, z_2) = F(w_2, w_1 | z_2, z_1) = F_n\left(\frac{z_2 w_2}{z_1}, \frac{z_1 w_1}{z_2} | z_1, z_2\right).$$

Thus the coefficients in the expansion $F_n(w_1, w_2 | z_1, z_2) = \sum_{0 \leq i, j \leq l} C_{i,j}(z_1, z_2) w_1^i w_2^j$ are rational functions in z_1, \dots, z_n obeying $C_{i,j}(z_1, z_2) = C_{j,i}(z_2, z_1) = \left(\frac{z_1}{z_2}\right)^{i-j} C_{j,i}(z_1, z_2)$. On the other hand from the explicit formula (17), one also finds that any $C_{i,j}(z_1, z_2)$ remains finite in the either limit $\frac{z_1}{z_2}, \frac{z_2}{z_i} \rightarrow \infty$ or $\frac{z_1}{z_2}, \frac{z_2}{z_i} \rightarrow 0$ for $i \geq 3$. It follows that $C_{i,j}(z_1, z_2) = 0$ unless $i = j$, hence

$$F_n(w_1, w_2, \dots, w_n | z_1, \dots, z_n) = F_n(1, w_1 w_2, w_3, \dots, w_n | z_1, \dots, z_n).$$

Moreover it is easily seen

$$F_n(1, w_1 w_2, w_3, \dots, w_n | z_1, z_2, \dots, z_n) = F_{n-1}(w_1 w_2, w_3, \dots, w_n | z_2, \dots, z_n).$$

Repeating this we reach $F_1(w_1 \dots w_n | z_n)$ giving the LHS of (17).

We note that the sum-to-1 (14) does not hold in the elliptic case. Remember that our local states are taken from $\mathcal{F} = \eta + \mathbb{Z}^{n+1}$ with a generic $\eta \in \mathbb{C}^{n+1}$. So we set $a = \eta + \tilde{a}$ with $\tilde{a} \in \mathbb{Z}^{n+1}$ etc in (4), and assume that it is valid also for $\tilde{a}, \tilde{b}, \tilde{c}, \tilde{d}$. It is easy to check

Proposition 1. *Assume $l \leq m$ and $|q| < 1$. Then the following equality holds:*

$$\lim_{\eta \rightarrow \infty} S_{l,m} \left(\begin{array}{c} \eta + \tilde{a} \quad \eta + \tilde{b} \\ \eta + \tilde{d} \quad \eta + \tilde{c} \end{array} \right)_{\text{trig}} = q^{\sum_{i < j} (\beta_i - \gamma_i) \gamma_j} \binom{m}{l}_q^{-1} \prod_{i=1}^{n+1} \binom{\beta_i}{\gamma_i}_q, \quad (18)$$

where the limit means $\eta_i - \eta_{i+1} \rightarrow \infty$ for all $1 \leq i \leq n$, and the RHS is zero unless $0 \leq \gamma_i \leq \beta_i, \forall i$.

The limit reduces the unrestricted trigonometric $A_n^{(1)}$ face model to the vertex model at a special value of the spectral parameter in the sense that the RHS of (18)| $_{q \rightarrow q^2}$ reproduces [8, eq.(23)] that was obtained as the special value of the quantum R matrix associated with the symmetric tensor representation of $U_q(A_n^{(1)})$.

3 Discussion

Since the weights $W_{l,m} \left(\begin{smallmatrix} a & b \\ c & d \end{smallmatrix} \middle| u \right)$ remain unchanged by shifting $a, b, c, d \in \tilde{\mathcal{P}}$ by $\text{const} \cdot (1, \dots, 1)$, we regard them as elements from $\mathcal{P} := \tilde{\mathcal{P}}/\mathbb{C}(1, \dots, 1)$ in the sequel. Given $l, m_1, \dots, m_M \in \mathbb{Z}_{\geq 1}$ and $u, w_1, \dots, w_M \in \mathbb{C}$, the transfer matrix $T_l(u) = T_l \left(u \middle| \begin{smallmatrix} m_1, \dots, m_M \\ w_1, \dots, w_M \end{smallmatrix} \right)$ of the unrestricted $A_n^{(1)}$ face model with periodic boundary condition is a linear map on the space of independent row configurations on length M row $\bigoplus \mathbb{C} |a^{(1)}, \dots, a^{(M)}\rangle$ where the sum is taken over $a^{(1)}, \dots, a^{(M)} \in \mathcal{P}$ such that $a^{(i+1)} - a^{(i)} \in B_{m_i}$ ($a^{(M+1)} = a^{(1)}$). Its action is specified as $T_l(u) |b^{(1)}, \dots, b^{(M)}\rangle = \sum_{a^{(1)}, \dots, a^{(M)}} T_l(u) \left(\begin{smallmatrix} a^{(1)}, \dots, a^{(M)} \\ b^{(1)}, \dots, b^{(M)} \end{smallmatrix} \middle| a^{(1)}, \dots, a^{(M)} \right)$ in terms of the matrix elements

$$T_l(u) \left(\begin{smallmatrix} a^{(1)}, \dots, a^{(M)} \\ b^{(1)}, \dots, b^{(M)} \end{smallmatrix} \right) = \prod_{i=1}^M W_{l, m_i} \left(\begin{smallmatrix} a^{(i)} & a^{(i+1)} \\ b^{(i)} & b^{(i+1)} \end{smallmatrix} \middle| u - w_i \right) \quad (a^{(M+1)} = a^{(1)}, b^{(M+1)} = b^{(1)}). \quad (19)$$

Theorem 1 tells that $S_l := T_l(u)_{u=w_1=\dots=w_M}$ has a simple factorized matrix elements. We write its elements as $S_{l, b^{(1)}, \dots, b^{(M)}}^{a^{(1)}, \dots, a^{(M)}}$. The star-triangle relation (7) implies the commutativity $[T_l(u), T_{l'}(u')] = [S_l, S_{l'}] = 0$.

Let us consider whether $X = T_l(u)$ or S_l admits an interpretation as a Markov matrix of a discrete time stochastic process. The related issue was treated in [3] for $n = 1$ and mainly when $\min(l, m_1, \dots, m_M) = 1$. One needs (i) sum-to-1 property $\sum_{a^{(1)}, \dots, a^{(M)}} X_{b^{(1)}, \dots, b^{(M)}}^{a^{(1)}, \dots, a^{(M)}} = 1$ and (ii) nonnegativity $\forall X_{b^{(1)}, \dots, b^{(M)}}^{a^{(1)}, \dots, a^{(M)}} \geq 0$. We concentrate on the trigonometric case in what follows. From Theorem 1 and the fact that $S_{l,m} \left(\begin{smallmatrix} a & b \\ c & d \end{smallmatrix} \right)_{\text{trig}}$ in (13) is independent of a , (i) indeed holds for S_l . On the other hand (13) also indicates that (ii) is not valid in general without confining the site variables in a certain range. A typical such prescription is *restriction* [4, 6, 5], where one takes $L = \ell + n + 1$ in (2) with some $\ell \in \mathbb{Z}_{\geq 1}$ and lets the site variables range over the finite set of level ℓ dominant integral weights $\{(L + a_{n+1} - a_1 - 1)A_0 + \sum_{i=1}^n (a_i - a_{i+1} - 1)A_i \mid L + a_{n+1} > a_1 > \dots > a_{n+1}, a_i - a_j \in \mathbb{Z}\}$. They are to obey a stronger adjacency condition [5, p546, (c-2)] than (3) which is actually the fusion rule of the WZW conformal field theory. (The formal limit $\ell \rightarrow \infty$ still works to restrict the site variables to the positive Weyl chamber and is called ‘‘classically restricted’’.) Then the star-triangle relation remains valid by virtue of nontrivial cancellation of unwanted terms. However, discarding the contribution to the sum (14) from those b not satisfying the adjacency condition spoils the sum-to-1 property. For example when $(n, l, m) = (2, 1, 2)$, $a = (2, 1, 0)$, $c = (4, 2, 0)$, $d = (3, 1, 0)$ and ℓ is sufficiently large, the unrestricted sum (14) consists of two terms $S_{l,m} \left(\begin{smallmatrix} a & b \\ d & c \end{smallmatrix} \right)_{\text{trig}} = \binom{2}{1}_q^{-1} \frac{(q^{-1}; q)_1}{(q^{-2}; q)_1}$ for $b = (4, 1, 0)$ and $S_{l,m} \left(\begin{smallmatrix} a & b' \\ d & c \end{smallmatrix} \right)_{\text{trig}} = \binom{2}{1}_q^{-1} \frac{(q^3; q)_1}{(q^2; q)_1}$ for $b' = (3, 2, 0)$ summing up to 1, but b' must be discarded in the restricted case since $a \stackrel{m=2}{\Rightarrow} b'$ [5, (c-2)] does not hold.

Thus we see that in order to satisfy (i) and (ii) simultaneously one needs to resort to a construction different from the restriction.

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References

1. G. E. Andrews, R. J. Baxter and P. J. Forrester, Eight vertex SOS model and generalized Rogers-Ramanujan-type identities, *J. Stat. Phys.* **35**, (1984) 193–266.
2. R. J. Baxter, *Exactly solved models in statistical mechanics*, Academic Press, London (1982).
3. A. Borodin, Symmetric elliptic functions, IRF models, and dynamic exclusion processes, arXiv:1701.05239.
4. E. Date, M. Jimbo, A. Kuniba, T. Miwa and M. Okado, Exactly solvable SOS models II: Proof of the star-triangle relation and combinatorial identities, *Adv. Stud. in Pure Math.* **16** (1988) 17-122.
5. M. Jimbo, A. Kuniba, T. Miwa and M. Okado, The $A_n^{(1)}$ face models, *Commun. Math. Phys.* **119** (1989) 543–565.
6. M. Jimbo, T. Miwa and M. Okado, Symmetric tensors of the $A_{n-1}^{(1)}$ family, *Algebraic Analysis*, **1** (1988) 253–266.
7. J. Kuan, An algebraic construction of duality functions for the stochastic $U_q(A_n^{(1)})$ vertex model and its degenerations, arXiv:1701.04468.
8. A. Kuniba, V. V. Mangazeev, S. Maruyama and M. Okado, Stochastic R matrix for $U_q(A_n^{(1)})$, *Nucl. Phys.* **B913** (2016) 248–277.