

Hurwitz theory, with (a) spin

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- Overview:**
- 1) (Spin) Hurwitz theory
 - 2) Fermion formalism & topological recursion
 - 3) Intersection theory on $\bar{\mathcal{M}}_{g,n}$
 - 4) Applications to Gromov-Witten theory

1) (SPIN) HURWITZ THEORY

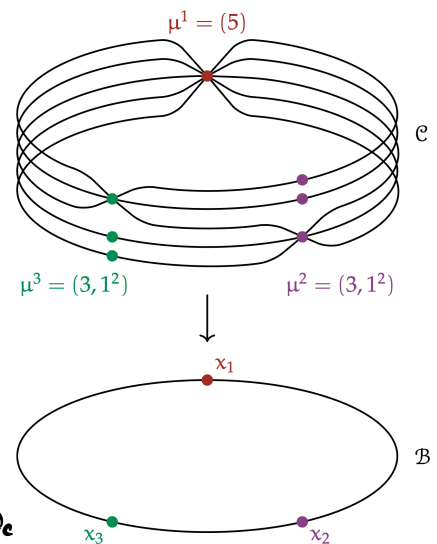
"(Spin) Hurwitz theory is the (signed) enumeration of branched covers of Riemann surfaces,

Defn. Fix a base $\mathcal{B} \ni x_1, \dots, x_k$. Fix $\mu^1, \dots, \mu^k + d$.

$$H_d(\mathcal{B}; \mu^1, \dots, \mu^k) := \sum_{[\mathcal{L}]} \frac{1}{|\text{Aut}(\mathcal{L})|} \quad \text{Hurwitz numbers}$$

iso classes of $f: \mathcal{C} \xrightarrow{d:1} \mathcal{B}$, \mathcal{C} unpt, cmtd, Riemann srfc

- ramified over x_i , ramification profile μ^i
- unramified everywhere else



- Questions:
- Properties?
 - How to compute?
 - What are they used for?

- Answers:
- Relation to rep. theory of S_d
 - Integrable hierarchies
 - Topological recursion
 - Intersection theory on $\bar{\mathcal{M}}_{g,n}$
 - Gromov-Witten theory of curves
 - Volumes of moduli spaces of holomorphic differentials
- $\left. \begin{array}{l} \text{• Relation to rep. theory of } S_d \\ \text{• Integrable hierarchies} \\ \text{• Topological recursion} \\ \text{• Intersection theory on } \bar{\mathcal{M}}_{g,n} \end{array} \right\} \mathcal{B} = \mathbb{P}^1$
 $\left. \begin{array}{l} \text{• Gromov-Witten theory of curves} \\ \text{• Volumes of moduli spaces of} \\ \text{holomorphic differentials} \end{array} \right\} \mathcal{B} = T$

Motivated by the geometry of the moduli space of holom. diffts, Eskin-Okounkov-Pandharipande introduced spin Hurwitz numbers.

Defn. A **spin structure** on B is a line bundle $\mathcal{S} \rightarrow B$ s.t. $\mathcal{S}^{\otimes 2} \cong \omega_B$. Define the **parity** $p(\mathcal{S}) \equiv h^0(B, \mathcal{S}) \pmod{2}$

↑ positive/negative

Expl: $\mu = (2m_1+1, \dots, 2m_k+1)$

- $\mathcal{O}(-1)$ is the only spin structr on P^1
- If $\mathcal{S} \rightarrow B$ is a spin structure and $C \xrightarrow{f} B$ a ramified cover with **ODD ramifications**, $\mathcal{S}_{f,B} = f^*\mathcal{S} \otimes \mathcal{O}(\frac{1}{2} \text{Ram})$ is a spin structure on C .

Defn. Fix a base $B \ni x_1, \dots, x_k$ and $\mathcal{S} \rightarrow B$ a spin structr. Fix $\mu^1, \dots, \mu^k + d$ **ODD partitions**.

$$H_d(B, \mathcal{S}; \mu^1, \dots, \mu^k) := \sum_{[\mathcal{L}]} \frac{(-1)^{p(\mathcal{S}_{f,B})}}{|Act(\mathcal{L})|}$$

Spin Hurwitz numbers

We have a monodromy map to the **spin symmetric group** :

$$0 \rightarrow \mathbb{Z}/2\mathbb{Z} \rightarrow \tilde{S}_d \rightarrow S_d \rightarrow 0$$

$\left\{ \begin{array}{l} [f: C \xrightarrow{d:1} B] \\ \text{positive or negative} \end{array} \right\} \xrightarrow{\sim} \left\{ \begin{array}{l} \text{repres of } S_d \\ \text{w/ or w/o a lift to } \tilde{S}_d \end{array} \right\}$

$f \mapsto \rho: \pi_1(B - \{\text{branch pts}\}, o) \rightarrow S_d$
 $\gamma \mapsto [p_i \mapsto \tilde{\gamma}_i(1)]$
 unique lift of γ starting at p_i

$f^{-1}(o) = \{p_1, \dots, p_d\}$
 $\lambda_1 > \lambda_2 > \dots$

FACTS: 1) $\{\text{irreps of } \tilde{S}_d\} \xrightarrow{\sim} \{\lambda + d \text{ strict partition}\}$

2) $\{\text{basis of the spin class algebra } \tilde{\Sigma}_d\} \xrightarrow{\sim} \{\mu + d \text{ odd partition}\}$

$$\tilde{f}_\mu(\lambda) = \text{spin central character} \sim \tilde{\chi}_\lambda(\mu)$$

Thm (**Character formula**, Eskin-Okounkov-Pandharipande '08, Gunningham '16). Disconnected spin Hurwitz numbers of $(B, \mathcal{S}) = (P^1, \mathcal{O}(-1))$ are given by

$$H_d^0(P^1, \mathcal{O}(-1); \mu^1, \dots, \mu^k) = 2^{\frac{\sum_i (e(\mu^i) - d)}{2} - d} \sum_{\substack{\lambda + d \\ \text{strict part}}} \left(\frac{\dim(\lambda)}{2^{PC(\lambda)/2} d!} \right)^2 \prod_{i=1}^k \tilde{f}_{\mu^i}(\lambda)$$

RESTRICTION: • $(B, \mathcal{D}) = (\mathbb{P}^1, \mathcal{O}(-1))$

- $\mu^1 = \mu$ generic **odd** partition of d
- $\mu^2 = \mu^3 = \dots = (r+1, 1, \dots, 1) + \text{lower order terms} =: \tilde{c}_r$ **r EVEN**

genus of the cover \uparrow

$$h_{g;\mu}^r := \frac{|\text{Act}(\mu)|}{b!} H_d(\mathbb{P}^1, \mathcal{O}(-1); \mu, \underbrace{\tilde{c}_r, \dots, \tilde{c}_r}_{b \text{ times}}) \quad b = \frac{2g-2+\ell(\mu)+d}{r}$$

$$= \# \left\{ \begin{array}{c} \text{C}_{g, \mathcal{D}} \text{ positive} \\ \downarrow \\ \mathbb{P}^1 \end{array} \right\} - \# \left\{ \begin{array}{c} \text{C}_{g, \mathcal{D}} \text{ negative} \\ \downarrow \\ \mathbb{P}^1 \end{array} \right\}$$

Character formula: $h_{g;\mu}^{o,r} = \frac{|\text{Act}(\mu)|}{b!} 2^{1-g-2d} \sum_{\substack{\lambda \vdash d \\ \text{strict part}}} \binom{\dim(\lambda)}{2^{P(\lambda)/2} d!} \tilde{f}_\mu(\lambda) \left(\frac{P_{r+1}(\lambda)}{r+1} \right)^b$

power sum
↓
 $P_{r+1}(\lambda)$



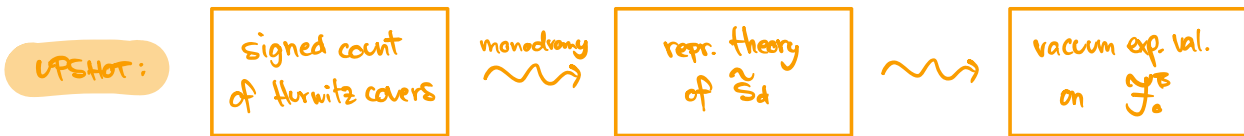
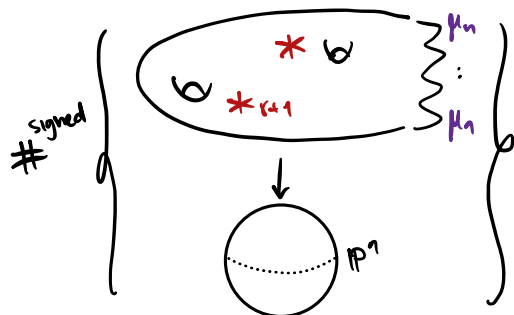
2) FERMION FORMALISM & TOPOLOGICAL RECURSION

- FACTS: $\exists \cdot (\mathcal{F}_0^B = \text{span} \{ |\lambda\rangle \mid \lambda \text{ strict part.} \}, \langle \cdot | \cdot \rangle, |0\rangle) = \text{Fock space of type B}$
- ↑ vect. space
 - ↑ pairing
 - ↑ distinguished element (vacuum)
- $\hat{D}_\infty \curvearrowright \mathcal{F}_0^B$ with explicit elements $J_{-\mu}^B$, $\mu \vdash d$ odd, F_{r+1}^B , r even, st.
- ↑ Lie algebra

$$J_{-\mu}^B |0\rangle = \sum_{\substack{\lambda \vdash d \\ \text{strict part.}}} \frac{\tilde{\chi}_\lambda(\mu)}{2^{P(\lambda)/2} \ell(\mu)} |\lambda\rangle \quad F_{r+1}^B |\lambda\rangle = P_{r+1}(\lambda) |\lambda\rangle$$

Prop (Spin Hfts as vacuum expectation values).

$$h_{g;\mu}^{o,r} = \frac{2^{1-g}}{b!} \langle 0 | e^{J_0^B} \left(\frac{F_{r+1}^B}{r+1} \right)^b \prod_{i=1}^n \frac{J_{-\mu_i}^B}{\mu_i} | 0 \rangle$$



ADVANTAGE:

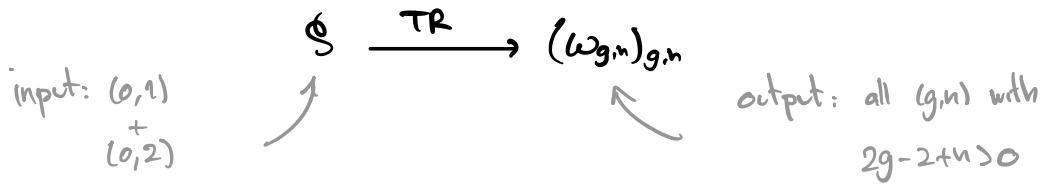
- We can use commutation relations of operators in $\hat{\mathcal{D}}_{\text{low}}$ to get closed formulae for low $(g, \ell(\mu))$
- The generating series of spin Hfts is a BKP τ -fnctn

Prop. For $(g, \ell(\mu)) = (0, 1)$ and $(0, 2)$, we get explicit formulae:

$$h_{0;\mu}^r = \frac{\mu^{\frac{\mu-1}{2}-2}}{\left(\frac{\mu-1}{r}\right)!}, \quad h_{0;\mu_1, \mu_2}^r = \frac{r}{\mu_1 + \mu_2} \frac{\lfloor \frac{\mu_1}{r} \rfloor! \lfloor \frac{\mu_2}{r} \rfloor!}{\lfloor \frac{\mu_1}{r} \rfloor! \lfloor \frac{\mu_2}{r} \rfloor!}$$

$\mu \in \mathbb{Z}_+^{\text{odd}}, r | \mu$ $\mu_1, \mu_2 \in \mathbb{Z}_+^{\text{odd}}, r | \mu_1 + \mu_2$

RECALL: given some initial data \mathcal{S} (a spectral curve $\mathcal{S} = (\Sigma, \pi, y: \Sigma \rightarrow \mathbb{C}, \mathcal{B})$), **topological recursion** is a recursive procedure that constructs a sequence of differentials $(\omega_{g,m})_{g,m}$ on the curve:



Thm (Topological recursion for spin H₁'s, conjectured in [GKL], proved by Alexandrov-Shadrin).

Spin H₁'s are computed by topological recursion on the spectral curve on \mathbb{P}^1

$$x(z) = \log(z) - z^r, \quad y(z) = z, \quad B(z_1, z_2) = \frac{1}{2} \left(\frac{1}{(z_1 - z_2)^2} + \frac{1}{(z_1 + z_2)^2} \right) dz_1 dz_2$$

by expanding the correlators near $e^{x(z_i)} = 0$:

$$\omega_{g,n}(z_1, \dots, z_n) = d_1 \cdots d_n \sum_{\substack{\mu \\ \text{odd part}}} h_{g;\mu}^r e^{\mu_1 x(z_1)} \cdots e^{\mu_n x(z_n)}$$



3) INTERSECTION THEORY on $\overline{\mathcal{M}}_{g,n}$

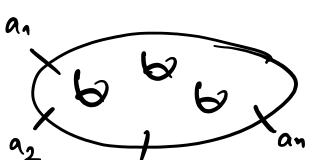
RECALL: for g,n st. $2g-2+n > 0$, the moduli space of stable curves

$$\overline{\mathcal{M}}_{g,n} := \left\{ (C, p_1, \dots, p_n) \mid C \text{ is a genus } g \text{ stable curve with } n \text{ smooth marked pts} \right\} / \sim$$

is a smooth cmet complx orbifold of $\dim_{\mathbb{C}} = 3g-3+n$. Thus, we can consider intersection numbers:

$$\int_{\overline{\mathcal{M}}_{g,n}} \alpha \in \mathbb{Q}, \quad \alpha \in H^{2(3g-3+n)}(\overline{\mathcal{M}}_{g,n}, \mathbb{Q})$$

EXPECTATION: every curve-counting problem can be expressed as an intersection number on the moduli space of curve:

$$N_{g; a_1, \dots, a_n} = \# \left\{ \begin{array}{l} \text{"structures" on} \\ \text{curve} \end{array} \right\} \stackrel{?}{=} \int_{\overline{\mathcal{M}}_{g,n}} \alpha_{a_1, \dots, a_n}$$


Question. What is the RHS for spin H^{tt}'s?

Thm (Spin ELSV formula). For $r=2$, spin H^{tt}s are given by double Hodge integrals:

$$h_{g; \mu_1, \dots, \mu_n}^{r=2} = 2^{4g-4+2n} \left(\prod_{i=1}^n \frac{\mu_i^{\frac{\mu_i-1}{2}}}{(\frac{\mu_i-1}{2}!)!} \right) \int_{\overline{\mathcal{M}}_{g,n}} \frac{\lambda(1) \lambda(-1/2)}{\prod_{i=1}^n (1 - \frac{\mu_i}{2} \psi_i)}$$

The formula is equivalent to topological recursion for spin H^{tt}s.

More generally, for arbitrary r (even), $\lambda(1) \lambda(-1/2)$ is substituted by a product of Witten's 2-spin class and Chiodo's class.

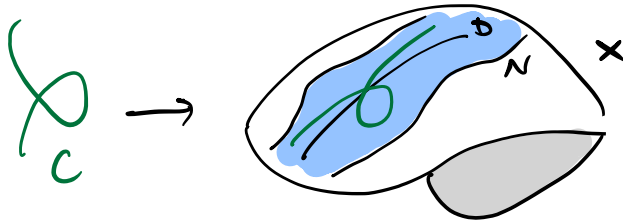
4) APPLICATIONS TO GROMOV-WITTEN THEORY (work in progress w/ Kromer, Lewański, Sawget)

Let X be a smooth variety, $\beta \in H_2(X, \mathbb{Z})$. GW invariants count curves in X of given genus and degree:

$$GW_g(X; \beta) = \# \left\{ C \xrightarrow{\mathbb{R}} X \mid \begin{array}{l} C \text{ is a curve of} \\ \text{genus } g \\ \mathbb{R}_* [C] = \beta \end{array} \right\}$$

If X is a Kähler surface of general type with a smooth canonical divisor D .

then GW invariants of degree dD , $d \in \mathbb{Z}_+$, localise around D :



Here N is the normal bundle.

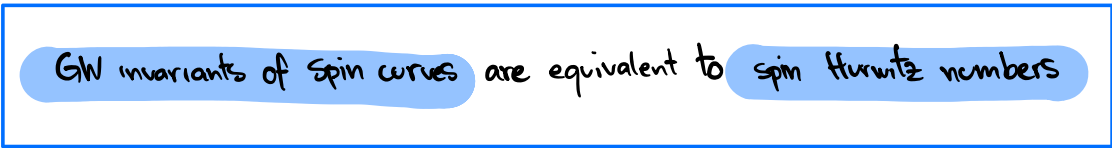
FACT: by the adjunction formula, N is a spin structure on D : $N^{\otimes 2} \cong \omega_D$.

Thm (Lee-Parker, Maulik-Pandharipande, Kiem-Li). $GW_g(X; dD)$ are determined by "local" invariants $GW_g^{\text{spin}}(D, N; d)$ that depend on the spin curve (D, N) .



Question. How to solve GW theory of spin curves?

Conjecture (Spin GW/H correspondence):



The proof (in progress) makes use of the spin ELSV formula as a starting point:

